

# **Operational Forecasting of Supercell Motion: Review and Case Studies Using Multiple Datasets**

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## ABSTRACT

The notion of anticipating supercell motion with multiple datasets in an operational setting is addressed. In addition, the most common propagation mechanisms that regulate both supercell and nonsupercell thunderstorm motion are reviewed. At minimum, supercell motion is governed by advection from the mean wind and propagation via dynamic vertical pressure effects. Therefore, one can use a hodograph to make predictions of supercell motion *before* thunderstorms develop, or before thunderstorms split into right- and left-moving components. This allows for better situational awareness and pathcasts of severe weather (relative to what occurs without *a priori* knowledge of supercell motion), especially during the early stages of a supercell's lifetime. There are several potential sources of wind data readily available across the United States, making it relatively easy to derive an ensemble of supercell motion estimates.

### 1. Introduction

Knowledge of supercell motion prior to thunderstorm formation is critical for successful short-term forecasting of the associated severe weather, which in turn is fundamental for emergency management and media preparations. For example, the forecast motion can be used to determine the storm-relative helicity (SRH), as well as the storm-relative flow at the low, middle, and upper levels of the supercell, which are important for evaluating tornadic potential and precipitation distribution (e.g., Rasmussen and Straka 1998; Thompson 1998; Rasmussen 2003). Storm motion forecasts may also be useful in assessing convective mode (e.g., LaDue 1998; Bluestein and Weisman 2000). Moreover, a reasonable forecast of supercell motion—relative to no knowledge of supercell motion—can lead to better pathcasts of hazardous weather in severe local warnings, especially during the initial stages of the supercell's lifetime (e.g., when a thunderstorm is beginning to split into right- and left-moving supercells). This is important since at least 90% of supercells are severe (e.g., Burgess and Lemon 1991) and most strong or violent tornadoes are produced by supercells (e.g., Moller et al. 1994). In addition, Bunkers (2002) found that 53 of 60 left-moving supercells produced severe hail (i.e., diameter  $\geq 1.9$  cm).

Modeling studies since the early 1980s (e.g., Weisman and Klemp 1986), field programs such as Verifications of the Origins of Rotation in Tornadoes Experiment (VORTEX; Rasmussen et al. 1994), and recent empirical studies (Bunkers et al. 2000; hereafter referred to as B2K) indicate that supercell motion can be anticipated prior to, and monitored during, severe weather operations. Despite the advances noted above, along with additional conference papers (e.g., Bunkers and Zeitler 2000; Edwards et al. 2002) and computer-based training (UCAR 1999), routine use of the B2K supercell motion forecast technique (discussed in section 2g) remains

limited to a few entities within the operational forecasting community [e.g., the Storm Prediction Center and a small percentage of the Weather Forecast Offices (WFOs)]. There may be a number of reasons for the lack of use, including simple unawareness, incomplete or inadequate training, lack of confidence in both the B2K technique and the general concept of forecasting supercell motion, or perceived lack of real-time data for the determination of vertical wind shear, supercell motion, and other derived parameters. The Sydney, NSW, Australia, hailstorm of April 14, 1999 (Bureau of Meteorology 1999) illustrates the perils in not understanding supercell processes.

This study emphasizes the mechanisms that influence supercell motion, and addresses the perceived lack of real-time data. A review of the most common mechanisms that control supercell, and thunderstorm, motion is presented in section 2 to provide relevant background information. Section 3 identifies near-real-time or real-time sources of wind profile data that can be used to evaluate vertical wind shear and predict supercell motion. Section 4 contains four case studies of supercell motion diagnosed or monitored by data sources listed in section 3. Section 5 presents conclusions and recommended actions for operational forecasters.

### 2. Mechanisms controlling supercell motion

A discussion of the mechanisms that control supercell motion necessarily includes those which affect the motion of nonsupercell thunderstorms. The two fundamental and distinctly different physical controls on storm motion are (1) *advection* and (2) *propagation*—the latter of which can be subdivided into several categories. In the current context, advection refers to the movement of a convective element with the mean flow as it entrains horizontal momentum into the updraft, whereas propagation refers to new convective initiation preferentially located relative to existing convection such that it has an overall effect on storm motion (i.e., the convection propagates through, not with, the mean flow). Propagation due to this new convective development requires the three primary ingredients for deep moist convection: moisture, instability, and upward motion. Based on a review of the literature, these two basic controls (advection and propagation) can be delineated as follows:

- (a) *advection* by the mean wind throughout a representative tropospheric layer,
- (b) *propagation* via dynamic vertical pressure gradients due to a rotating updraft (germane to supercells only),
- (c) *propagation* via convective development along a thunderstorm's outflow,
- (d) *propagation* via convective development along a boundary layer convergence feature,
- (e) *propagation* via storm mergers and interactions, and

- (f) *propagation* via orographic effects such as upslope flow, lee-side convergence, and an elevated heat source.

Other thunderstorm propagation mechanisms exist [e.g., propagation due to gravity waves; Carbone et al. (2002)], but those listed above are believed to represent the most significant influences on localized thunderstorm motion.

Storm motion is typically determined by tracking a prominent feature in radar reflectivity or velocity data (e.g., the storm centroid, the mesocyclone, etc.). Tracking the storm centroid gives the impression that a storm is an *object*, when in reality, the storm represents a *process* that is strongly affected by parcel ascent (the growing part of the storm) and descent (the dissipating part of the storm). Therefore, a storm continually changes due to these processes, and cannot be considered a solid object (Hitschfeld 1960; Doswell 1985, p. 52). The propagation components discussed in the present study largely represent storm processes of ascent.

Items (a) and (b) above are assumed to be the most important determinants of supercell motion (see B2K and section 2g). The remaining four items (and especially d–f), while at times very important, cannot be effectively anticipated with a hodograph on a consistent basis, and instead must be inferred from observational data such as surface, satellite, radar, and topography. However, it is noted that the general direction of (c) has potential predictability with a hodograph based on the orientation of the low-level vertical wind shear vector (Rotunno et al. 1988). Knowledge of these various thunderstorm/supercell propagation mechanisms (in addition to advection) is especially important when attempting to make an accurate forecast of supercell motion. Since there is still some disagreement regarding the processes that influence supercell motion [e.g., see exchange by Klimowski and Bunkers (2002) and Weaver et al. (2002a,b)], these mechanisms are discussed in the following paragraphs.

#### a) *Advection*

The most fundamental mechanism controlling the motion of deep moist convection is advection by the mean wind throughout a representative tropospheric layer (e.g., Brooks 1946; Byers and Braham 1949), which is analogous to (but not the same as) flow in a stream. Hitschfeld (1960) described this process as conservation of horizontal momentum in rising or descending parcels of air, which is tempered by entrainment. Utilizing radar data, these early thunderstorm studies revealed the motion of *discrete nonsupercell storms* was highly correlated with advection by the mean cloud-bearing wind [also see summary by Chappell (1986, pp. 293–294)]. There have been many variations on what tropospheric layer is appropriate to calculate the mean wind, but some commonly used methods have included the mandatory sounding levels up to 300 or 200 hPa (e.g., Newton and

Fankhauser 1964; Maddox 1976) or the layer from the surface to 6 km (~20000 ft) (e.g., Byers and Braham 1949; Weisman and Klemp 1986; B2K)—the top of which is often less than the maximum cloud height. In support of this shallower layer, Wilhelmson and Klemp (1978) noted that the motion of modeled thunderstorms did not change appreciably when the wind speeds above 6 km were substantially increased. By way of contrast, Wilson and Meigenhardt (1997) found that the 2–4 km layer gave the best results in predicting convective cell motion in Florida. More recently, Ramsay and Doswell (2004) suggested a deeper layer for advection of supercells (e.g., surface to 8 km). Obviously the “correct” advection depth is dependent on factors such as the height of the storm (e.g., low-topped supercells are advected over a shallower layer compared to tall storms), whether the storm’s inflow is rooted in a near-surface layer or an elevated mixed layer (e.g., nocturnal convection north of a warm front would not be affected by near-surface winds), or perhaps the storm’s age (e.g., younger storms may be advected over a shallower layer relative to older storms). Advection may also be regionally dependent, thus “scaling” the mean wind depth is potentially useful (Thompson et al. 2004).

A frequently overlooked consideration in the calculation of the mean wind is whether or not pressure-weighting (or similarly, density-weighting) is employed. Taking the average of the mandatory pressure-level winds implicitly involves pressure-weighting because of the greater concentration of mandatory levels in the lower half of a sounding. In this way, it is assumed that advection dominates in the lower atmosphere—relative to the upper atmosphere—due to the exponential decrease of pressure (density) with height. Pressure-weighting also has been used in shallower atmospheric layers (e.g., Weisman and Klemp 1986), although it is unclear if this is necessary based on the lack of statistical studies relating nonsupercell thunderstorm motion to the mean wind. Clearly, this is an area that would benefit from additional research.

The effect of the mean wind on thunderstorm motion increases as the wind speed increases. When the mean wind, and hence advection, is weak ( $< 10 \text{ m s}^{-1}$ ), other propagation mechanisms (described below) can have a significant impact on storm motion, but when the mean wind is strong ( $> 20 \text{ m s}^{-1}$ ), it tends to dominate over the propagation mechanisms. In the absence of advection (i.e., the mean wind is near zero), the only other control on thunderstorm motion is propagation. The most relevant propagation mechanisms are described throughout the remainder of this section.

#### b) *Shear-induced propagation*

The feature that sets supercells apart from other thunderstorms is their persistent, rotating updraft in the midlevels of the storm (Fig. 1; Doswell and Burgess 1993; Moller et al. 1994). Although a multicellular

structure can be imposed upon a supercell, a rotating updraft lasting at least tens of minutes in the lower- to mid-levels of a thunderstorm defines it as a supercell. This rotation results from the interaction of the updraft with the vertically sheared environment, whereby low-level horizontal vorticity ( $\omega_H$ ; Figs. 1a & 2a) is tilted into the vertical and becomes spatially associated with the updraft (Figs. 1b & 2b,c). The rotating updraft, in turn, produces a localized region of lowered pressure in the midlevels of the storm, which creates an upward-directed pressure gradient force (Klemp 1987; Fig. 1). With respect to the three primary ingredients for deep moist convection (i.e., moisture, instability, and upward motion), this pressure gradient force provides enhanced upward motion on a preferred storm flank to initiate new convective development. Numerical modeling studies have shown that this dynamic interaction can contribute around 50% of the total updraft strength (Weisman and Rotunno 2000; see their Fig. 13). This effect produces a horizontal updraft–shear propagation (USP) component that is perpendicular to the shear vector and anti-parallel (parallel) to the horizontal vorticity vector for the right-moving (left moving) supercell (Fig. 2b,c; also see Weisman and Klemp 1986). It is this basic attribute of supercells that largely explains the often disparate motion of right- and left-moving supercells (Davies-Jones 2002).

Davies-Jones (2002) further subdivided the shear effects into linear and nonlinear components. The nonlinear component is described above, and the linear component is a function of hodograph curvature which is maximized for circular hodographs. It is beyond the scope of this paper to discuss the details of these differences, but both effects lead to propagation that is perpendicular to the shear vector.

In general, USP manifests itself through continuous (or perhaps quasi-continuous) movement of a supercell (as opposed to discrete movement), although this is a function of the temporal and spatial resolution of the radar data. This propagation may also depend on the strength of the vertical wind shear [i.e., stronger shear may lead to a larger propagation component (Bunkers and Zeitler 2000)]. Propagation due to this updraft–shear interaction becomes increasingly important as the mean wind speed decreases, the vertical wind shear increases, and also for atypical hodograph orientations (see B2K). When the USP opposes the mean wind, slow-moving supercells can result.

#### c) *Gust-front propagation*

As a thunderstorm matures and develops precipitation, it produces a relatively cold outflow at the surface (Fujita 1959). The leading edge of this outflow acts as a lifting mechanism, providing enhanced upward motion relative to the existing storm. New convective cells may form due to the convergence of moist and unstable air along the leading edge of this outflow, which

can become a flanking line if positioned on the side of the storm (Lemon 1976). Daughter or feeder cells along this outflow or flanking line may interact with the main thunderstorm to affect its overall motion and intensity (e.g., Fig. 3; Lemon 1976; Browning 1977; Weaver and Nelson 1982). Browning (1977) defined daughter cells as convective elements which grow and become the dominant thunderstorm cell while the original cell decays; this process often leads to discrete/broken thunderstorm propagation. Browning defined feeder cells as convective elements which merge with and intensify the main thunderstorm but do not become dominant; this process typically leads to quasi-discrete/quasi-continuous propagation (e.g., Fig. 3). When discrete propagation from daughter cells occurs along a supercell’s gust front, it can result in a multicell-supercell hybrid as described in Weaver and Nelson (1982) and Foote and Frank (1983).

In terms of forecasting, the gust front or cold pool strength depends upon cloud base height, sub-cloud and mid-tropospheric relative humidity, and mid- to upper-level storm relative winds. These parameters provide guidance for the downdraft strength, which is subsequently affected by evaporation and loading by hydrometeors near the updraft. The strength of the cold pool generally increases when the midlevel shear is weak ( $\sim 0.003 \text{ s}^{-1}$ ) and dry air exists in midlevels and/or below cloud base. Development of new cells usually occurs in the direction of the low-level shear vector (Rotunno et al. 1988), so a hodograph may provide qualitative guidance for the direction of the gust-front propagation, but more quantitative results have been elusive (e.g., Bunkers and Zeitler 2000). This propagation mechanism also depends upon the moisture availability and the degree of instability in advance of the gust front.

This “gust-front” type of propagation, whether from daughter or feeder cells, can slow a system down when the flanking line is located on the upwind side of the thunderstorm, but it can accelerate the storm system when the gust front moves ahead of the storm [e.g., analogous to Corfidi’s (2003) discussion of forward-propagating mesoscale convective systems (MCSs)]. It is not known how often this phenomenon occurs with supercells, but Klemp (1987) noted that this type of propagation may play a secondary role for supercells when compared to the USP discussed in section 2b. Since low-precipitation (LP) supercells are less likely to have a significant cold pool, relative to non-LP supercells, gust-front propagation is least likely with these supercells.

#### d) *Boundary layer convergence features*

A common method of propagation for thunderstorm complexes is due to convective development along boundary layer convergence features such as fronts, drylines, moisture/ instability axes, and outflow boundaries, which often leads to discrete thunderstorm propagation (e.g., Newton and Fankhauser 1964; Weaver 1979; Magsig et al. 1998). This propagation mechanism

is not to be confused with gust-front propagation discussed in section 2c, although the two are interrelated when a gust front is interacting with a boundary layer convergence zone. This mechanism is driven by convergence of moist and unstable air that is often enhanced near boundaries (e.g., Weaver 1979; Maddox et al. 1980), which results in an increased probability of the initiation of deep moist convection (Wilson and Schreiber 1986). Furthermore, this process is often aided by the transport of moist and unstable air in a low-level jet, especially for MCSs (Corfidi et al. 1996).

Propagation due to boundary layer convergence features is different from (b) and (c) above since the existing thunderstorm is not enhancing upward motion; rather, the enhanced upward motion is derived from the convergence zone. Therefore, if a supercell is close enough to a boundary layer convergence zone, it may preferentially “move” along or toward this feature due to the development of new convective updrafts proximate to this favorable environment. Moreover, supercells or multicell-supercell hybrid storms may propagate against the mean flow in these low-level convergence zones due to mergers with daughter or feeder cells, resulting in little overall movement of the storm system (e.g., Fig. 4; Weaver 1979; Wakimoto et al. 2004). A classic example of this “boundary layer convergence” type of propagation occurred in the Jarrell, TX, tornadic event (27 May 1997), where new thunderstorm development was consistently along a pre-existing wind-shift boundary—leading to discrete thunderstorm propagation to the southwest as opposed to being a continuously propagating distinct supercell (Magsig et al. 1998).

It is important to note the “storm motion” can be computed in two different ways for this type of scenario. Referring to Fig. 4, the initial supercell has a northeastern motion of  $\mathbf{V}_1$ . At some later time, a new cell develops near the boundary layer convergence zone and merges with the supercell. This discrete propagation produces a system motion of  $\mathbf{V}_2 < \mathbf{V}_1$ , but the original (and weakening) cell still has a motion much closer to  $\mathbf{V}_1$ . After this process has completed, the motion of the entire system is much slower than that of the initial supercell via discrete propagation (i.e.,  $\mathbf{V}_3 < \mathbf{V}_2 < \mathbf{V}_1$ ), but individual cell motion may still be similar to  $\mathbf{V}_1$  (i.e., the old decaying supercell). If care is not taken in making the calculations, one can incorrectly assign the “system” motion to the “cell” motion by assuming that the original cell and the new cell are the same, when in fact cells are moving faster, or through, the system. In this example, the system (cell) motion results from discrete (continuous) propagation. This example is similar to the Superior, NE, supercell described by Wakimoto et al. (2004).

Boundary layer convergence may be especially important in environments characterized by large buoyancy and relatively weak shear [i.e., bulk Richardson number (BRN)  $> 50$ ], such as was the case for Jarrell, TX. This scenario also would often be characterized by a weak

mean wind ( $< 10 \text{ m s}^{-1}$ ), thus thunderstorms would be more likely to stay close to boundary layer convergence zones for a longer period of time than when the mean wind is strong ( $> 20 \text{ m s}^{-1}$ ). Newton and Fankhauser (1964) attempted to explain the deviant motions (relative to the mean wind) as being related to water-budget constraints, such that the largest storms—relative to the small storms—moved farthest to the right of the mean wind, thereby “intercepting” a larger volume of water vapor. Using a numerical cloud model, Atkins et al. (1999) studied the interaction of simulated supercells with preexisting boundary layer convergence zones and found that the effect on storm motion was about  $5 \text{ m s}^{-1}$ . When propagation due to boundary layer convergence dominates thunderstorm motion, locally heavy rainfall becomes increasingly likely.

#### e) Storm mergers and interactions

Very few studies have addressed the effects of storm mergers on subsequent storm motion, although considerable anecdotal evidence exists for its occurrence. In general, this occurs when there is an intersection of the paths of two thunderstorms whereby hydrometeors are redistributed between the storms, and their updrafts and downdrafts interact. For example, downdrafts may merge to initiate new convection via bridging (Westcott 1994), or the downdraft from one storm may enhance the downdraft of another storm, causing it to accelerate. There are times when storm mergers and interactions may lead to the intensification of existing convection (Lemon 1976), but there are other times when convection may be affected negatively by a merger (Westcott 1984). The effect of mergers on storm motion has not been investigated as extensively as the effects of the other propagation mechanisms discussed herein. However, in a study of bow-echo evolution across the United States, Klimowski et al. (2004) showed that when thunderstorms merge to become a bow echo, the subsequent motion of the bow echo was often dictated by the most dominant and aggressive cell prior to merger, which was usually the cell that initiated the merger (e.g., Fig. 5). This often resulted in an acceleration of the storm.

Cell mergers become increasingly likely when there are (1) numerous thunderstorms, (2) differing storm motions, and (3) strong linear synoptic forcing. Thunderstorms are most numerous when moisture, instability, and upward motion are relatively large, and convective inhibition and vertical wind shear are relatively weak. Differing storm motions can occur for a number of reasons: splitting thunderstorms and ordinary storms, short and tall storms being advected by different atmospheric layers, and discrete propagation due to gust fronts, boundary layer convergence features, and orography (discussed below). Finally, strong linear synoptic forcing acts to concentrate thunderstorms along a common feature (e.g., front or dryline), which then promotes adjacent cell interactions (e.g., Bluestein and

Weisman 2000). Storm mergers can also occur in conjunction with gust fronts (e.g., Fig. 3) and boundary layer convergence zones (e.g., Fig. 4), making it difficult to separate the two mechanisms, and at times leading to an “anchoring” of convection.

*f) Orographic effects*

Orography can also significantly influence supercell and nonsupercell thunderstorm motion. Elevated terrain can provide enhanced mesoscale convergence zones in at least three ways: (1) upslope flow, (2) lee-side convergence, and (3) an elevated heat source (Fig. 6; Kuo and Orville 1973; Banta 1990). New cells may continually develop in these favored regions as old ones advect away, leading to discrete thunderstorm propagation such that the system motion is near zero—effectively “anchoring” the thunderstorm system to the elevated terrain (e.g., Akaeda et al. 1995). As a result, this effect is very similar to that of boundary layer convergence zones discussed in section 2d, but the forcing mechanisms are clearly distinct. Nearby storms that are not affected by mountainous terrain may have a substantial motion relative to those being influenced by the orography. This effect appears to be most important when the mean wind is relatively weak ( $< 10 \text{ m s}^{-1}$ ) so that storms are not rapidly advected away from the source of mesoscale convergence.

Orography can also have other effects on storms, such as causing their demise in a downslope flow region or modulating their intensity by disrupting the inflow. Although this does not affect storm motion directly, it does affect storm evolution, which is an important forecast consideration.

*g) Predicting supercell motion*

B2K used 260 hodographs from supercell environments to develop a method to predict supercell motion, and they also reviewed methods that were applicable to the prediction of supercell motion. Their method assumes that advection and USP are the dominant mechanisms controlling supercell motion [effects a–b above; Fig. 2]; these two effects can be easily computed with a hodograph. [A tutorial on hodograph interpretation can be found in Doswell (1991) and UCAR (2003).] Using a hodograph and the B2K method, supercell motion can be predicted as follows (shown schematically in Fig. 7):

- (1) plot a representative mean wind, which may be derived from the surface to 6 km (B2K), the surface to 8 km (Ramsay and Doswell (2004), or a layer above the surface for elevated supercells (e.g., Thompson et al. 2004);
- (2) draw a shear vector that extends from the boundary layer (BL) to 5.5–6 km;

- (3) draw a line that both passes through the mean wind and is orthogonal to the shear vector (this represents the USP component; refer to section 2b); and
- (4) plot the right-moving (left moving) supercell motion  $7\text{--}8 \text{ m s}^{-1}$  from the mean wind, and along the orthogonal line to the right (left) of the shear vector.

In addition to plotting the forecast supercell motion on a hodograph, the B2K method can be applied to numerical model data, and therefore storm motion vectors can be overlaid on radar imagery, as is demonstrated in section 4c (also see Klimowski and Bunkers 2002).

There are times when several of the six items discussed above (advection plus five propagation mechanisms) might play a role in supercell motion, and this reduces the effectiveness of B2K’s technique. Moreover, poor supercell motion estimates—from any method—may occur due to (1) use of an unrepresentative sounding, (2) use of an inappropriate mean wind layer, and (3) varying deviations from the mean wind due to USP, which may be dependent upon the strength of the vertical wind shear or to thermodynamic considerations. However, B2K showed a  $1\text{--}2 \text{ m s}^{-1}$  improvement in mean absolute error over other methods available to predict supercell motion (e.g., 30 degrees to the right of the mean wind and 70 percent of the mean wind speed). More recently, Edwards et al. (2002) and Ramsay and Doswell (2004) found the B2K method to be statistically superior to other supercell motion forecasting schemes, and Edwards et al. (2004) applied the B2K method to a dataset of 32 left-moving supercells. Therefore, this technique will be used in section 4 to show how supercell motion can be anticipated operationally.

### 3. Sources of vertical wind information

*a) Radiosonde observations*

A scan of WFO area forecast discussions (AFDs) shows that radiosonde observations (RAOBs) are still the most popular source of data for tropospheric wind information. This is understandable from the standpoint that RAOBs have been the traditional data source spanning the advent of modern meteorology after World War II. In addition, RAOBs have the positive attributes of (i) in-situ data, (ii) concurrent thermodynamic and wind data, and (iii) well-known and minimized equipment and acquisition errors. Unfortunately, RAOBs have two substantial limitations, especially with respect to forecasting and monitoring severe convection: (i) poor spatial resolution [69 sites in the Continental United States (CONUS) (Peterson and Durre 2004)], and (ii) poor temporal resolution (observations typically at 12-h intervals). Many individual storms, and even entire mesoscale convective systems (MCSs), can initiate, mature, and dissipate without being sampled by the

RAOB network. Forecasters attempt to remedy these deficiencies by modifying RAOBs to represent the current or forecast (temporal), and/or nearby (spatial) environment. Doswell (1991) discussed some of the pitfalls in these modifications, and Brooks et al. (1994) discussed the notion of proximity soundings. Overall, RAOBs are the best source of data if spatially and temporally near convection; however, in many instances this is not the case.

*b) 404 MHz vertical wind profilers*

The NOAA Profiler Network (NPN; formerly Wind Profiler Demonstration Network) provides a remotely sensed source of vertical wind information, primarily across the Great Plains (NOAA 1994). The primary advantages of the NPN are high temporal resolution (averaged observations at least hourly), and a vertical resolution of 250 m. These attributes are especially useful when monitoring the low-level jet across the central United States during convective situations. However, disadvantages of the NPN include only four sites outside the Great Plains (three in Alaska, and one in central New York—the majority are between 30 and 45 °N and 87 and 108 °W), precipitation attenuation, and contamination from biologic sources such as bird movements. NPN data have been found to be quantitatively consistent with the accuracy and reliability of RAOBs (NOAA 1994). Further information on the NPN can be found at: <http://www.profiler.noaa.gov/jsp/index.jsp>.

*c) Doppler radar vertical wind profiles*

The Weather Surveillance Radar-1988 Doppler (WSR-88D) vertical wind profiles (VWPs) are a remotely sensed source of data with similar temporal resolution to the NPN. In contrast to the NPN, the WSR-88D VWPs provide coverage across the CONUS, with the best spatial resolution over the southern and eastern CONUS, and the poorest over the western CONUS. VWP data are collocated with many CONUS RAOBs, which allows for comparison studies near the 0000 UTC and 1200 UTC RAOB launches, but this collocation also reduces the potential of sampling a larger geographical region if they were not collocated. The VWP data share many of the same limitations as the NPN due to precipitation attenuation and biologic contamination. Qualitatively, the VWP are of lesser quality than both RAOBs and the NPN (Don Burgess, personal communication, 2003). An additional disadvantage of the WSR-88D VWPs is the lack of data above the boundary layer prior to convection (due to a lack of scatterers) (Klazura and Imy 1993; see their Table 1).

*d) Aircraft Communication and Addressing and Reporting System (ACARS)*

Aircraft Communication Addressing and Reporting System (ACARS) observations provide a growing source of wind profile data. The advantages of ACARS data include in-situ, fast-response sensors, high vertical resolution, and on some aircraft, concurrent temperature and dew point observations. Spatial and temporal resolution is good near hub airports (e.g., ATL, DFW, LAX, ORD), but poor over the bulk of the CONUS except at altitudes above 7.5 km (~25,000 ft). However, new initiatives such as the Tropospheric Airborne Meteorological Data Reporting System (TAMDAR; Daniels 2002) offer the prospects of much higher spatial and temporal resolution below 7.5 km. Free data access is restricted to airlines, NOAA, and research groups, thereby limiting the utility and applications of ACARS data. Further details can be found at: <http://acweb.fsl.noaa.gov>.

*e) Model analysis/forecast soundings*

Another growing source of wind profile data are model analyses and forecasts. In fact, many small operating units such as WFOs and universities regularly run local models allowing for customization of resolution, domain, and physics. In general, model analyses are good, especially for pre-convective environment. Studies by Thompson et al. (2003), and others (see <http://maps.fsl.noaa.gov>), have shown fair to good agreement between model analysis profiles and nearby RAOBs. Excellent horizontal, vertical, and temporal resolution can be tuned to provide data in, or close, to the convective area of interest, reducing the time required to obtain a “proximity” sounding dataset for studying tornadic vs. nontornadic supercells (e.g., Thompson et al. 2003). However, disadvantages include errors introduced in the analysis or forecast process, which can lead to errors in analysis or forecast fields.

*f) Data sources summary*

In general, observed data are preferred over model analyses or forecasts. Similarly, in-situ data are preferred over remotely sensed observations, due to assumptions made in the remote sensing retrieval process that may render misleading data. Finally, close temporal and spatial proximity usually represents the near-storm environment better than observations further away. Sometimes, a mix of the various data is required to obtain a reasonable estimate of the environment in proximity to a given weather phenomenon (as the example in section 4a illustrates). Table 1 provides a subjective summary of these data types.

#### **4. Example cases**

In the following four examples, knowledge of supercell motion outlined in section 2 is used with the various datasets described in section 3 to show how

supercell motion can be effectively anticipated in an operational setting.

a) *Southeast Texas—May 30, 1999*

The Fort Bend County supercell of 30 May 1999 was one of several which formed as part of a "northwest flow" event, referring to winds from 270 to 360 degrees at or above 3000 m in response to an upper-level longwave ridge upstream and trough downstream of the convective area (e.g., Johns 1982). These events have atypical hodographs as defined by B2K, resulting in values of SRH and other storm-relative parameters to be unrepresentative of supercell potential when the storm motion is estimated with non-Galilean invariant methods. [As noted in B2K, Galilean invariant methods maintain the same storm motion forecast, relative to a given vertical wind shear profile, no matter what the ground-relative wind profile looks like. However, non-Galilean invariant methods can give different storm motion forecasts for the same vertical wind shear profile but with differing ground-relative winds.] The supercell produced two reports of 1.9 cm hail, one F1 tornado (which destroyed a barn and numerous trees and power poles), and a flash flood ("waist-deep" water in one subdivision). This was a relatively ordinary severe weather episode, but is still useful to illustrate the concepts of forecasting supercell motion operationally. It also highlights the importance of viewing supercell motion from the vertical wind shear perspective (Galilean invariant), and not the mean wind perspective (non-Galilean invariant).

Tracking the evolution and subsequent storm motion of the Fort Bend supercell was complicated by convection along its flanking line, as well as a second storm to its southeast (Fig. 8). Although the Fort Bend supercell remained identifiable throughout its lifetime, the convection along the flanking line periodically merged with the main storm (in the form of feeder cells), and apparently accelerated it to the south-southwest (as discussed in the next paragraph). Moreover, the second storm (Fig. 8) eventually merged with the Fort Bend supercell on its downshear (eastern) flank, although its affect on storm motion is unclear. Finally, this convection was occurring within a northeast-southwest axis of relatively moist and unstable air, which may have further aided storm propagation through boundary layer convergence since the storm-relative inflow was from the south-southwest. In summary, the Fort Bend supercell appeared to be influenced by advection and four of the five supercell propagation mechanisms discussed in section 2—orography being the only factor deemed unimportant.

The mean supercell motion for the Fort Bend supercell was from 18 degrees at  $7 \text{ m s}^{-1}$ . Local forecaster vernacular for this type of movement is a "southwest-moving supercell"—implying some special class of supercell. In reality, the supercell was moving to the right of the mean shear vector. This motion is the same as a

typical upper-right quadrant hodograph supercell; except that the shear vector was rotated roughly 90 degrees clockwise from typical orientations by the mean northwest flow. Figs. 9 through 12 show hodographs derived from four different sources: (i) the 0000 UTC 5/31/1999 Corpus Christi (CRP), TX, RAOB, (ii) the 0000 UTC 5/31/1999 Lake Charles, LA, (LCH) RAOB, (iii) the 2300 UTC 5/30/1999 MAPS analysis sounding nearest to Houston Hobby Airport (KHOU), and (iv) the 0014 UTC 5/31/1999 Houston/Galveston WSR-88D (KHGX) VWP. The 0000 UTC CRP hodograph produced the smallest predicted motion error at  $2.1 \text{ m s}^{-1}$ —although most of the data sources would have provided a reasonable estimate of storm motion (i.e., 2–5  $\text{m s}^{-1}$  errors), especially considering the complicating factors of gust-front propagation and storm mergers. Non-Galilean invariant methods, such as those based on a percentage of the mean wind speed and an angular deviation to the right, would have produced a supercell motion forecast that was to the southeast, and in error by as much as 7–8  $\text{m s}^{-1}$ . In general, the supercell moved faster—to the south and west—than predicted by the B2K method (e.g., Figs. 11 & 12). This is consistent with the merging feeder cells observed along the southwestern leading edge of the storm (Fig. 8), which would accelerate development along this flank of the supercell.

The values of 0-6 km total shear ( $U_s$ ), 0-6 km bulk shear ( $U_b$ ), and 0-3 km SRH were calculated for the wind profile sources listed above, plus three other plausible sources (Table 2).  $U_s$  was calculated by summing the shear segments across each 0.5-km sublayer from 0 to 6 km, and  $U_b$  was calculated by determining the vector difference between the surface and 6-km winds. Calculated  $U_s$  ranged from  $27.4 \text{ m s}^{-1}$  to  $48.9 \text{ m s}^{-1}$ ,  $U_b$  ranged from  $5.7 \text{ m s}^{-1}$  to  $21.6 \text{ m s}^{-1}$ , and SRH ranged from  $15 \text{ m}^2 \text{ s}^{-2}$  to  $220 \text{ m}^2 \text{ s}^{-2}$ .  $U_s$  is greater than or equal to  $U_b$  by definition, since hodograph curvature is neglected in the calculation of  $U_b$ .

Overall,  $U_s$  was the least variable parameter (in terms of proportionality between the largest to smallest values)—roughly 71%.  $U_b$  had a roughly 280% difference, which was due to a wide variation in the 6-km wind. However, if the 0000 UTC LCH RAOB is discarded (because of convective contamination in the mid-to-upper levels), and the 1200 UTC RAOBs are omitted because of their temporal unrepresentativeness, then  $U_b$  ( $U_s$ ) only varied by 38% (71%). SRH ranged an order of magnitude between largest and smallest, and when the smallest value of SRH was omitted, the values still varied by 120%. These results agree with Markowski et al. (1998), B2K, and Weisman and Rotunno (2000), that  $U_s$  and  $U_b$  are more consistent predictors of supercell potential than SRH (assuming initiation of deep moist convection).

Earlier points about the data sources are also evident. The 0014 UTC 5/31/99 KHGX VWP was the closest non-model data source to the near-storm environment, and more strongly indicated supercell potential than the other

sources. The 2300 UTC 5/30/99 MAPS and RUC analysis sounding statistics show that relatively small differences in  $U_s$  (or  $U_b$ ) can still be associated with SRH differing by an order of magnitude between them (Table 2). [MAPS is the development version of the RUC analysis/model run at the Forecast Systems Laboratory (FSL), whereas the RUC is the operational version run at the National Centers for Environmental Prediction (NCEP).] This example demonstrates that even minor changes to models can have significant impacts for assessment of the near storm environment.

*b) Central Texas—March 26, 2000*

The central Texas supercells of March 26, 2000 were isolated left-moving (LM) and right-moving (RM) components of a split from an initial thunderstorm near the Granger (KGRK), TX, WSR-88D (Fig. 13). The LM supercell produced 2.2 to 6.9 cm hail, wind gusts to  $31 \text{ m s}^{-1}$ , and wind damage to mobile homes and roofs. The RM supercell produced an F0 tornado four miles north of Seguin, TX, 4.4 to 6.3-cm hail, and widespread wind damage to windows, roofs, cars, and power lines.

Fig. 14 shows the hodograph derived from the 0100 UTC 27 March 2000 MAPS analysis sounding closest to KGRK. Unlike the Fort Bend supercell, these supercells were isolated, and did not appear to be affected in a measurable way by gust-front propagation, boundary layer convergence, storm mergers, or orography. Note the reasonable prediction of the motion of the LM ( $V_{LM}$ ) and RM ( $V_{RM}$ ) supercells from the method of B2K—compared to the observed supercell motions ( $V_{obs}$ ) from 2223 UTC 26 March 2000 (shortly after the supercell split) to 0100 UTC 27 March 2000 (errors  $1.9\text{-}2.4 \text{ m s}^{-1}$ ). This case represents a situation where no traditional wind data were readily available, but model analyses provided good insight into the shear environment, resulting in a proper estimate of storm motions for the LM and RM supercells.

This case also illustrates how situational awareness can be improved simply by anticipating the motion of supercells, especially at far ranges from the radar in cases where data may be limited or missing (e.g., Maddox et al. 2002; their Fig. 1). For example, if radar velocity data were missing for this event, but the operational forecaster anticipated the tracks of the right- and left-moving supercells, there would be a smaller chance of being surprised when the supercells began to evolve, identification of supercells would be more straightforward, and severe storm warnings could potentially be improved.

*c) South-central South Dakota—August 3, 2001*

On the afternoon of 3 August 2001, a single right-moving supercell (i.e., a thunderstorm with a counterclockwise-rotating updraft) was observed over south central South Dakota. Severe weather consisted of

two hail reports (1.9 cm and 2.5 cm), and a wind gust to  $27 \text{ m s}^{-1}$ . This supercell was highly unusual in that it moved toward the northwest, and to the *left* of the mean wind (discussed below). The synoptic setting contained a midlevel ridge with weak northwesterly flow, but moderate southerly flow was prevalent in the lower atmosphere.

The initial thunderstorm formed around 1915 UTC, and supercell characteristics became evident by 2000 UTC. The lifetime of this lone supercell thunderstorm (when it possessed rotation) was about 90 minutes, and some mesocyclone alarms were indicated by the National Severe Storms Laboratory (NSSL) mesocyclone detection algorithm (Stumpf et al. 1998). The supercell moved northwest at  $8 \text{ m s}^{-1}$  (Fig. 15), and the parent thunderstorm dissipated by 2200 UTC.

The hodograph derived from the wind profiler at Merriman, NE, revealed that the supercell moved to the *left* of the mean wind, but to the *right* of the vertical wind shear (Fig. 16). Note that the B2K method had a storm motion error of  $5.4 \text{ m s}^{-1}$  (compare  $V_{obs}$  with  $V_{RM-fcst}$ ); however, the B2K method provided useful guidance of the anomalous motion by indicating a westward movement, which was to the right (left) of the shear vector (mean wind). This information would greatly reduce the chance of being “surprised” by anomalous supercell motion.

Referring back to the propagation mechanisms discussed in section 2, gust-front propagation and boundary layer convergence were of similar importance in modulating storm motion when compared to USP. First, the 0-6 km bulk shear was  $13 \text{ m s}^{-1}$ , which is on the low end for supercell occurrence [e.g., see Fig. 2 in Bunkers (2002)]. As a result, one might expect USP not to be as significant as when the bulk shear is much stronger. Second, a sequence of radar images revealed a gust front moving to the northwest, and away, from the supercell thunderstorm. The gust front may have caused the supercell to accelerate toward the northwest as the storm tried to “keep up” with this lifting mechanism. Indeed, the thunderstorm dissipated after the gust front was 20 km ahead of it. Finally, the supercell occurred along a gradient of moisture (Fig. 17), which might have been a source of enhanced instability. As noted in section 2d, this can lead to new convective development. In support of this, there was at least one period of discrete propagation, toward the northwest, early in the supercell’s lifetime.

In summary, the supercell moved toward the northwest as a result of USP, gust-front propagation, and boundary layer convergence. These latter two external forcing mechanisms can become important when the wind shear is weak (as in this case). This unusual motion to the left of the mean wind can be anticipated by viewing supercell motion from a vertical wind shear perspective. Nearby wind profiler data were useful for assessing the vertical wind profile and estimating storm motion. If a forecaster is utilizing the various datasets discussed in

section 3 to anticipate supercell motion before thunderstorms develop, they should not be caught off guard by anomalous motion of supercells.

*d) South-central South Dakota—June 30, 2003*

During the late afternoon of June 30, 2003, ordinary thunderstorms over South Dakota eventually gave way to one dominant long-lived supercell which lasted five hours. The supercell initiated in south-central South Dakota, and traveled into north-central Nebraska before dissipating. There were ten severe hail reports, ranging from 1.9–7.0 cm (five were from golfball to baseball size). This event presented a warning challenge as storm motion changed abruptly when the thunderstorm transitioned into a supercell, making the initial warning decision difficult with respect to what area would be affected.

At 2247 UTC no severe storms or supercells were occurring in south-central South Dakota, but ordinary nonsevere thunderstorms were moving east-northeast at 4–5 m s<sup>-1</sup> (Fig. 18a). Between 2247 UTC and 2315 UTC, an ordinary storm rapidly developed into a supercell and commenced moving southward at 4–5 m s<sup>-1</sup> (Fig. 18b)—a change in direction of 100–120 degrees to the right. This change in direction was well forecast by the Rapid Update Cycle (RUC) model using the B2K supercell motion forecasting method (Fig. 19a). As the supercell traveled into Nebraska, it displayed a south-southwest motion, slightly farther to the right of the motion indicated from the Merriman, NE, wind profiler (Fig. 19b).

In contrast to the previous example, the 0–6 km bulk shear was over twice as large in this case (around 30 m s<sup>-1</sup>), suggesting a more significant USP component, which may explain the stronger rightward deviation with time (Fig. 19a,b). None of the other propagation mechanisms appeared to be playing a significant role in this case since discrete propagation was not evident, no mergers occurred, and orography was not a factor.

In summary, this last case illustrates the usefulness of overlaying the mean wind and supercell motion vectors on radar images when making warning decisions (also see Klimowski and Bunkers 2002). The result can be improved short-term forecasts and warnings of severe weather associated with supercells, especially during the initial stages of a supercell's lifetime.

## 5. Conclusions and recommendations

Forecasting and monitoring supercell motion is critical to effective severe weather operations. This understanding begins with a proper conceptual model of the factors that influence supercell motion, which have been reviewed herein (refer to section 2). Armed with this knowledge, operational forecasters can use all available datasets to anticipate supercell motion. RAOBs are too spatially and temporally coarse to provide accurate wind profiles for estimating supercell motion in over half

of all events. Fortunately, NPN wind profiles, WSR-88D VWP, ACARS, and model analysis and forecast profiles can serve as surrogates for improved estimates of the near-storm environment, resulting in better anticipation and forecasts of supercell potential and motion. However, each source has advantages and disadvantages that can render the acquired data either invaluable or nearly useless. The case studies from 30 May 1999, 26 March 2000, 3 August 2001, and 30 June 2003 are ordinary examples of how varied data sources can provide differing levels of accuracy for anticipating supercells and forecasting their motion.

The case studies presented herein only describe but a few ways in which supercell motion can be anticipated operationally. In general, one starts with a baseline prediction using a hodograph (or plan view display) and assumes advection and USP are dominant. This prediction can then be modified contingent upon the anticipation of the other propagation mechanisms, if they are deemed to be significant. Although prediction of supercell motion remains inexact, clearly the potential exists to make substantial improvements in storm motion forecasts when considering advection and the most common propagation mechanisms.

Despite the advances in supercell theory, observing systems, and operational modeling, the severe weather operations meteorologist is still faced with applying a preponderance of evidence in the selection of the most appropriate data source(s), especially for forecasts prior to storm development. Knowledge of the mechanisms that control supercell motion, along with the various data sources, parameter robustness, and preferred method for calculating storm motion, will provide the best results for severe weather operations.

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## TABLES

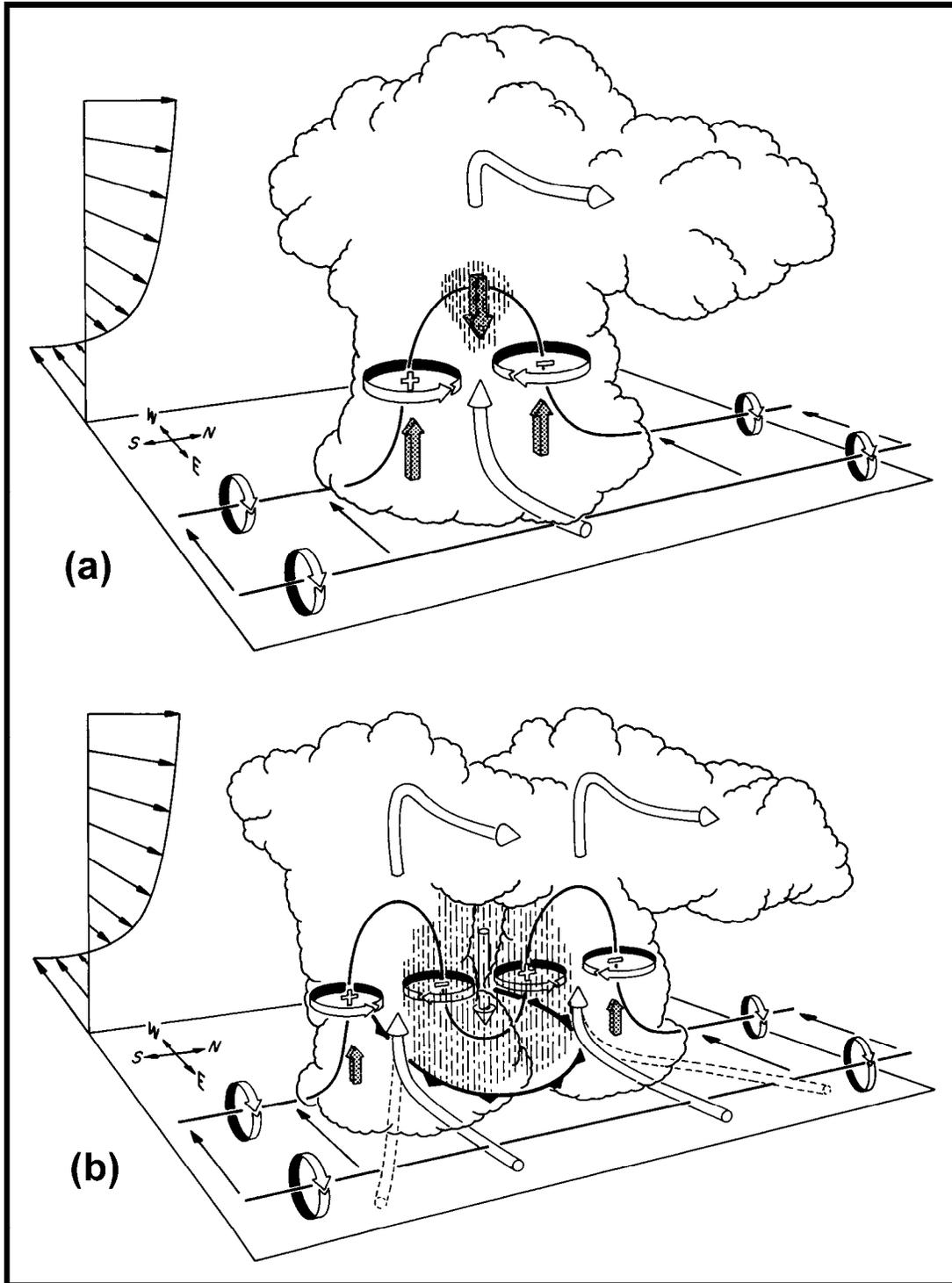
**Table 1.** Subjective summary of various aspects of the wind profile data discussed in section 3.

Data Source	Spatial/Temporal Resolution	Wind Error Bounds
RAOBs	~200 km/12 hr	+/- 1.5 m s <sup>-1</sup>
Wind Profilers	~200 km (central Plains)/6 min.	+/- 2.0 m s <sup>-1</sup>
WSR-88D VWP	~150 km/5-6 min.	+/- 3.0 m s <sup>-1</sup>
ACARS	~100 km (near hubs)/~1 hr	+/- 1.0 m s <sup>-1</sup>
Model Analyses/Forecasts	~1000 km (elsewhere)/~1 hr	+/- 2.0 m s <sup>-1</sup>
	~10 km/1 hr	+/- 2.0 m s <sup>-1</sup>

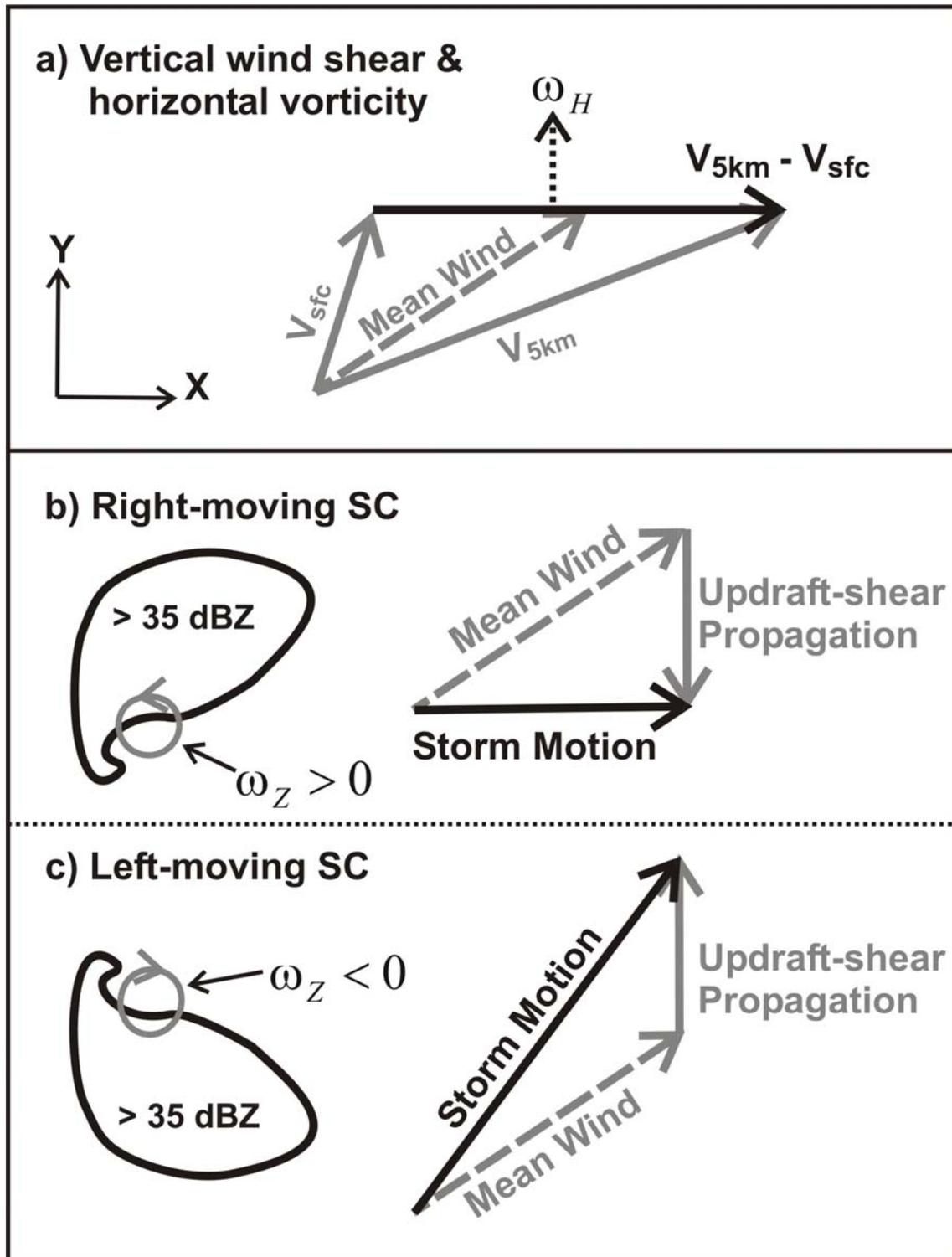
**Table 2.** 0-6 km total shear ( $U_s$ ), 0-6 km bulk shear ( $U_b$ ), and 0-3 km storm relative helicity (SRH) for potential storm environment data sources.

	0-6 km Total Shear (m s <sup>-1</sup> )	0-6 km Bulk Shear (m s <sup>-1</sup> )	0-3 km SRH (m <sup>2</sup> s <sup>-2</sup> )
12 UTC, LCH	31.0	5.7	105
12 UTC, CRP	46.8	12.8	120
00 UTC, LCH	35.0	5.7	183
00 UTC, CRP	27.4	21.6	100
23 UTC, MAPS	29.4	15.7	15
23 UTC, RUC	33.7	18.5	122
0014 UTC, KHGX	48.9	21.5	220

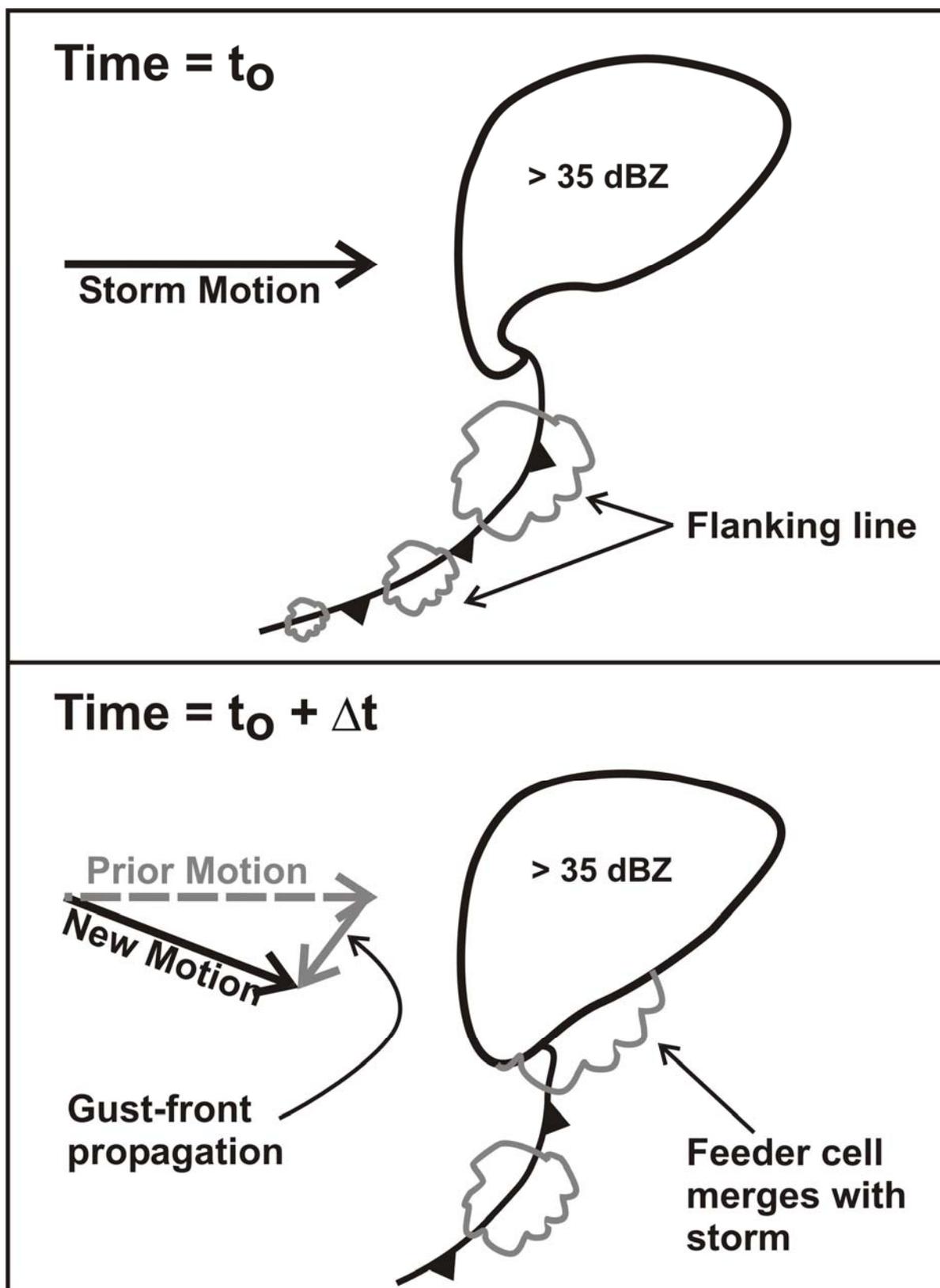
## FIGURES



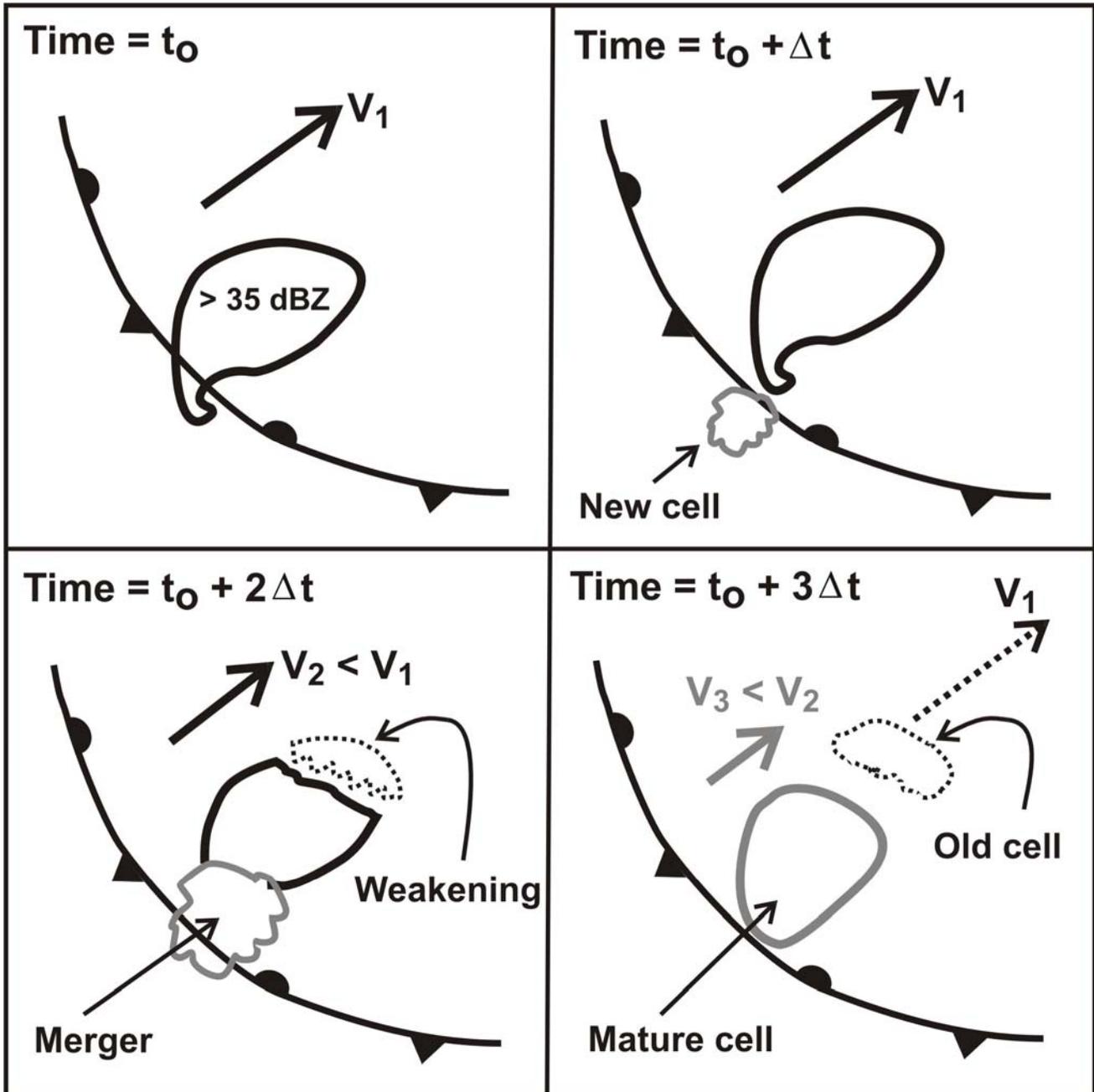
**Fig. 1.** Schematic illustrating the process of thunderstorm splitting in an environment of westerly vertical wind shear. In (a) horizontal vorticity,  $\omega_H$ , is tilted into the updraft to produce positive (negative) vertical vorticity,  $\omega_Z$ , on the storm's southern (northern) flank, which then locally enhances the updraft. In (b) the storm splits into cyclonically ( $\omega_Z > 0$ ) and anticyclonically ( $\omega_Z < 0$ ) rotating supercells as a downdraft forms between the two updraft maxima, with the two separate rotating updrafts propagating normal to the vertical wind shear. Reprinted from Klemp (1987), with permission, from the *Annual Review of Fluid Mechanics*, Volume 19 ©1987 by Annual Reviews, [www.annualreviews.org](http://www.annualreviews.org).



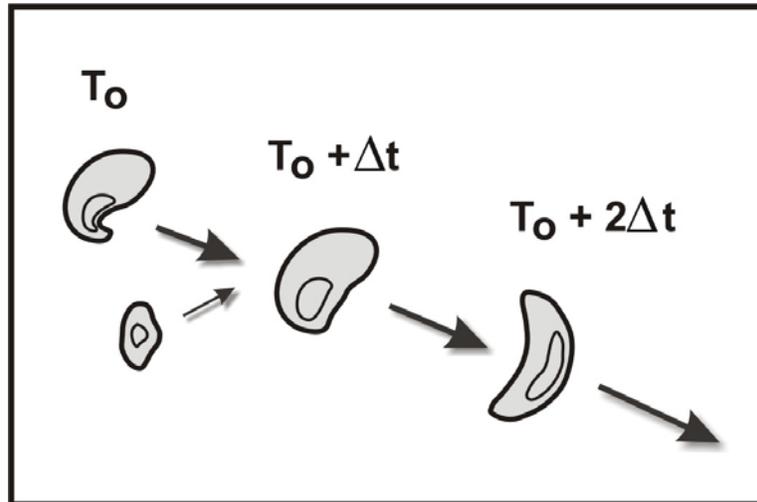
**Fig. 2.** Schematic illustrating the relationship between the vertical wind shear and supercell propagation. In (a) the 0-5 km vertical wind shear is given by  $V_{5km} - V_{sfc}$ , the mean wind is depicted with the grey dashed arrow, and the horizontal vorticity,  $\omega_H$ , which is perpendicular to the vertical wind shear, is represented with the dotted arrow. In (b) and (c) the mean wind is depicted with the grey dashed arrows, the propagation due to the rotating updraft is given by the grey arrows, and the resultant storm motion is given by the black arrows. Grey circles associated with the supercells represent the updraft circulation on the storm flank, and  $\omega_z$  is used to signify the vertical vorticity associated with the supercells which has arisen from the tilting of  $\omega_H$ .



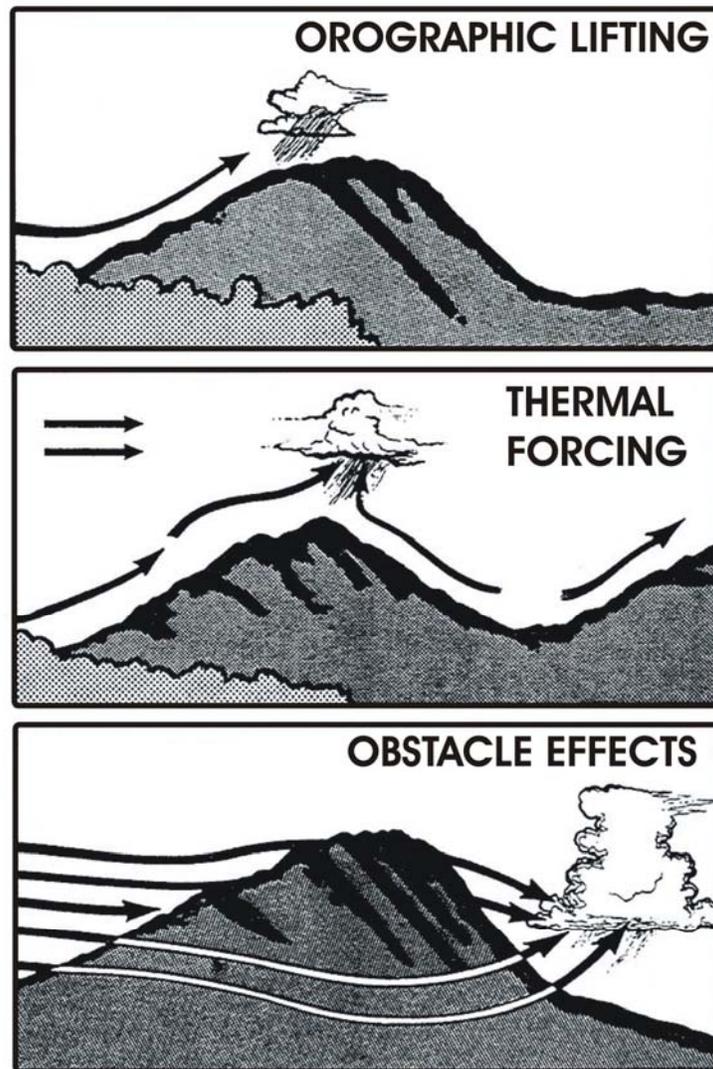
**Fig. 3.** Schematic illustrating the effect new cell development along a storm's gust front, or flanking line, can have on storm motion. At time  $t_0$  no cells are merging with the main storm, although convective development is noted along the storm's outflow. At time  $t_0 + \Delta t$  a convective cell has merged with the main storm, which may produce a propagation component in the direction of the merger.



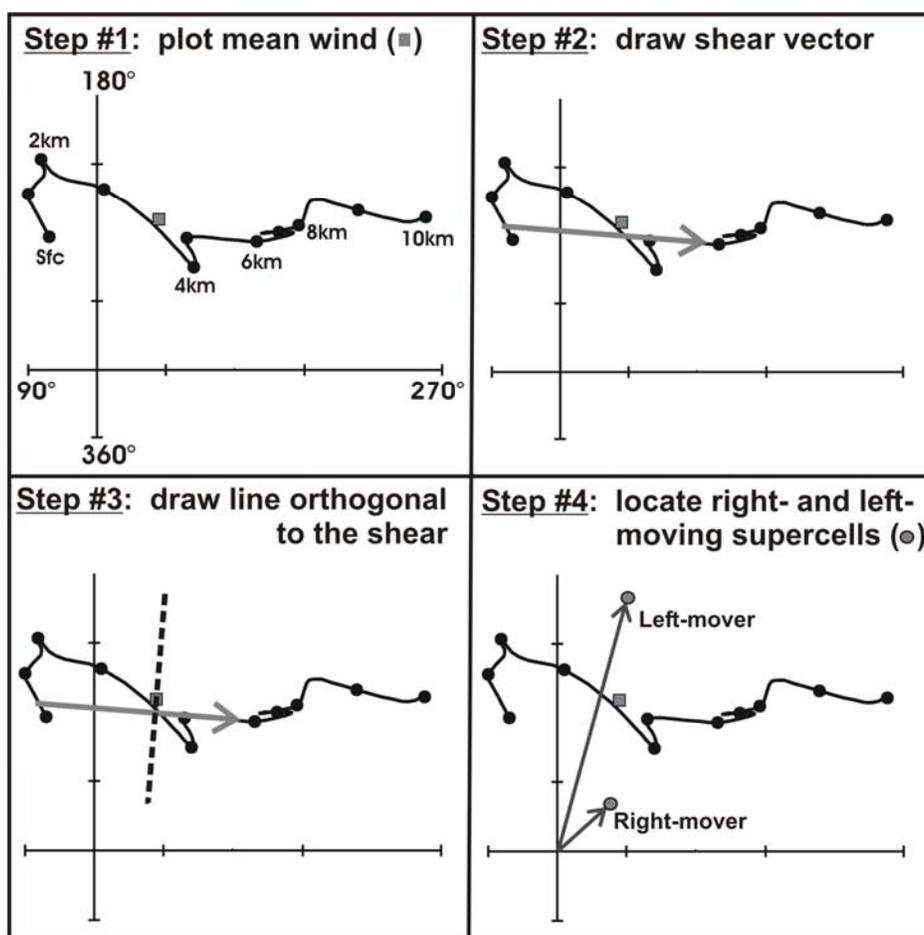
**Fig. 4.** Schematic illustrating the effect new cell development along a boundary layer convergence zone (in this case a stationary front) can have on storm motion. At time  $t_0$  a supercell is moving with velocity  $V_1$ . At times  $t_0 + \Delta t$  and  $t_0 + 2\Delta t$  a new convective cell forms along the convergence zone and merges with the initial supercell—slowing the overall system motion to  $V_2 < V_1$ . Finally, at time  $t_0 + 3\Delta t$  a new mature cell has evolved with an effective system motion of  $V_3 < V_2$ , while the old decaying supercell continues moving away to the northeast at about  $V_1$ .



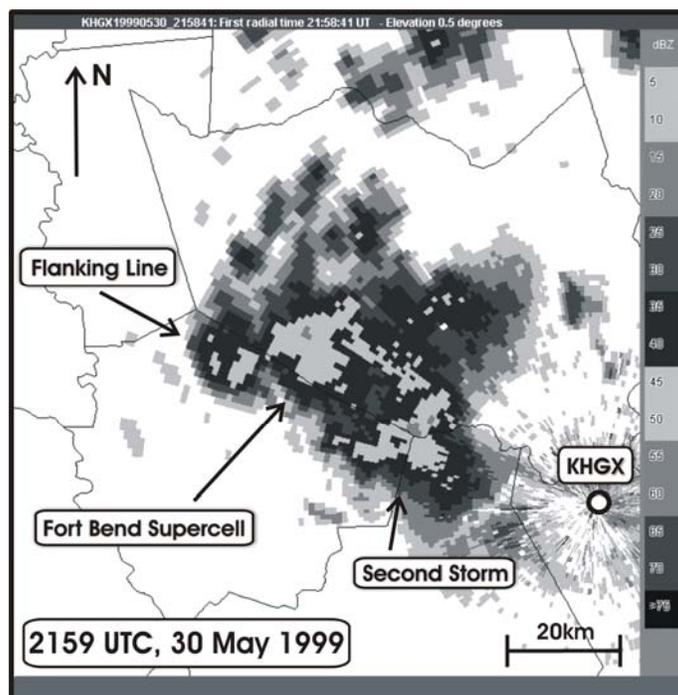
**Fig. 5.** Schematic illustrating one possible effect of storm mergers on storm motion. Adapted from Klimowski et al. (2004).



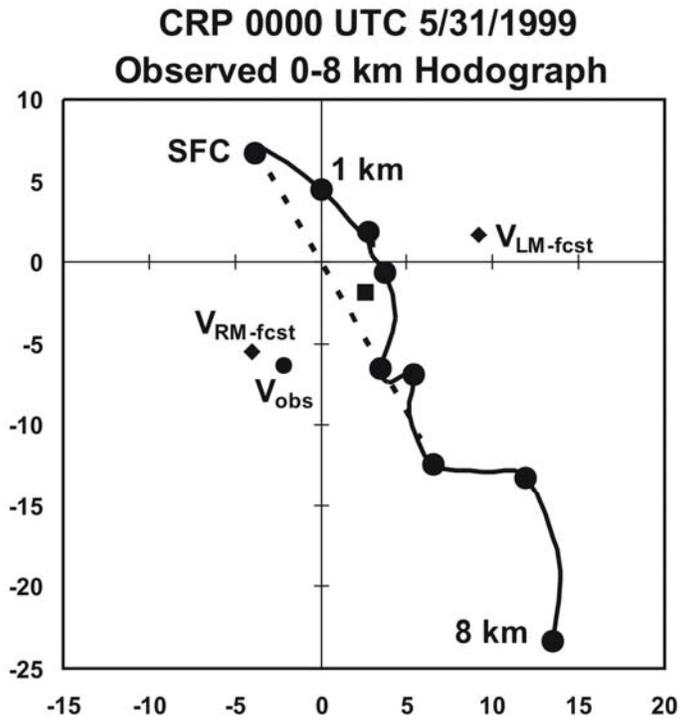
**Fig. 6.** Schematic illustrating various effects of orography on convective development, which can subsequently affect storm motion. Reproduced from Banta (1990).



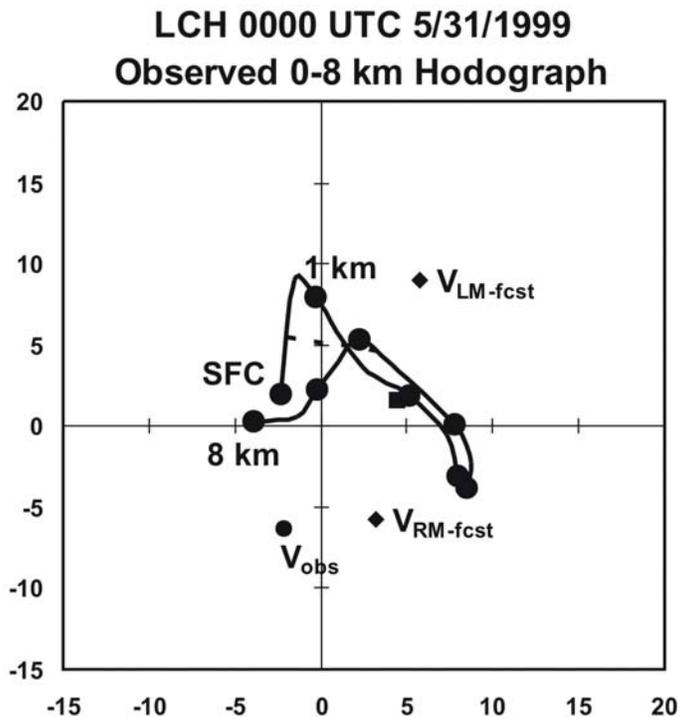
**Fig. 7.** Schematic illustrating the steps in plotting the predicted supercell motion using a hodograph and the method from Bunkers et al. (2000), which is discussed in section 2g.



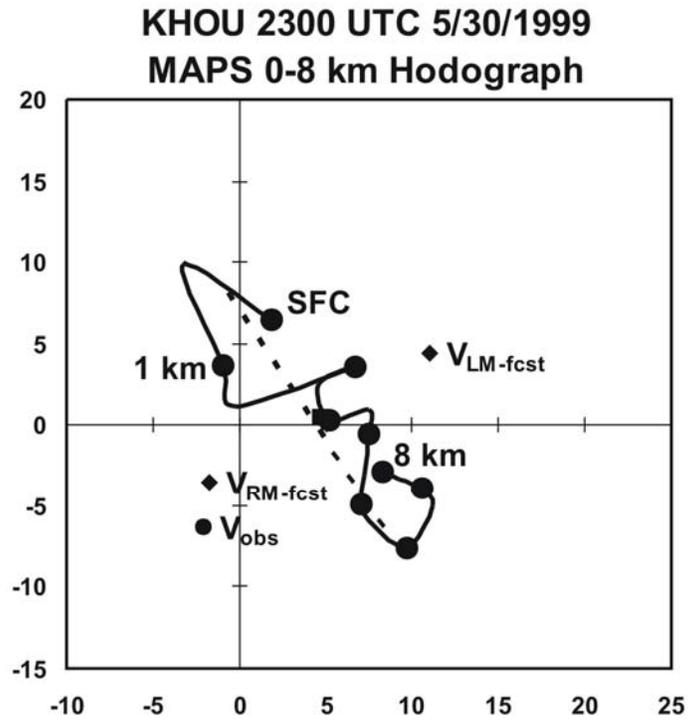
**Fig. 8.** KHGX WSR-88D 0.5° reflectivity image at 2159 UTC 30 May 1999. The Fort Bend supercell is discussed in section 4a. The Java-IRAS software was used to display the radar data (Preignitz 1995).



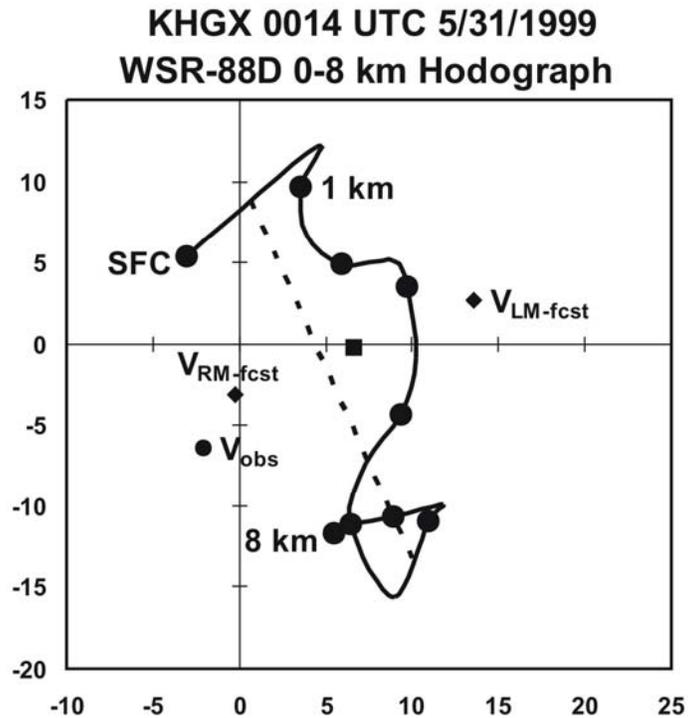
**Fig. 9.** The 0000 UTC 31 May 1999 CRP RAOB hodograph.  $V_{RM-fcst}$  is predicted right-moving supercell motion from the method in Bunkers et al. (2000).  $V_{LM-fcst}$  is the predicted left-moving supercell motion.  $V_{OBS}$  is the observed supercell motion. The dashed line represents the surface to 6 km shear. The square is the surface to 6 km mean wind. All units are  $m s^{-1}$ .



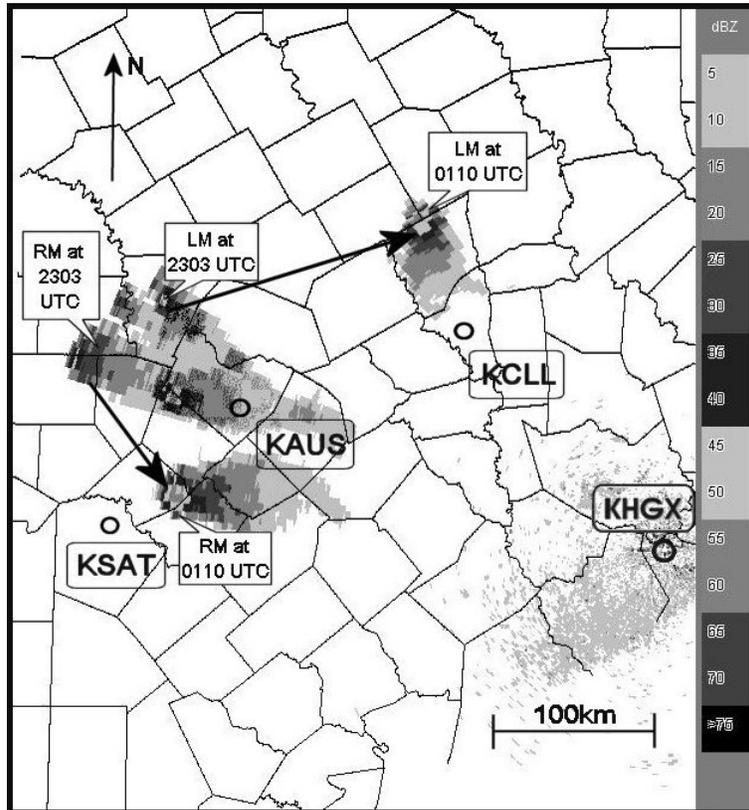
**Fig. 10.** Same as Fig. 9, except for the 0000 UTC 31 May 1999 LCH RAOB hodograph.



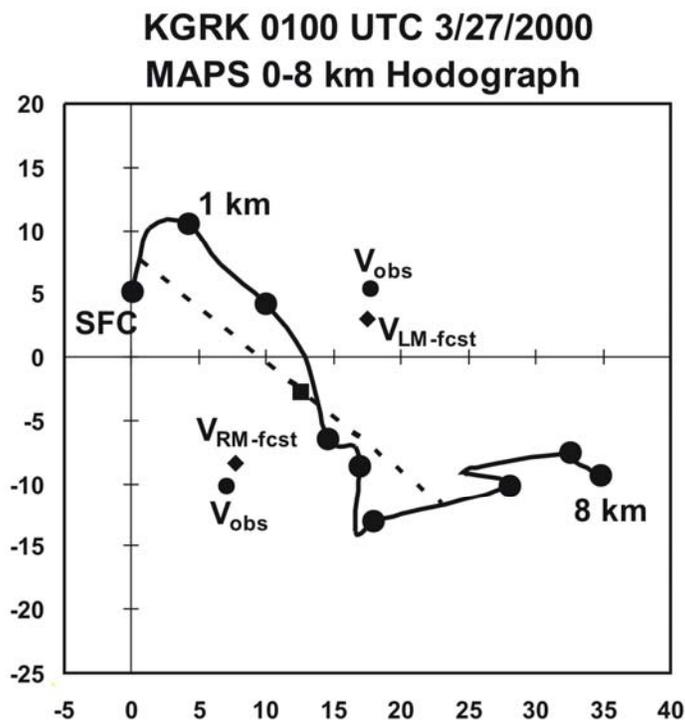
**Fig. 11.** Same as Fig. 9, except for the 2300 UTC 30 May 1999 MAPS (KHOU) hodograph.



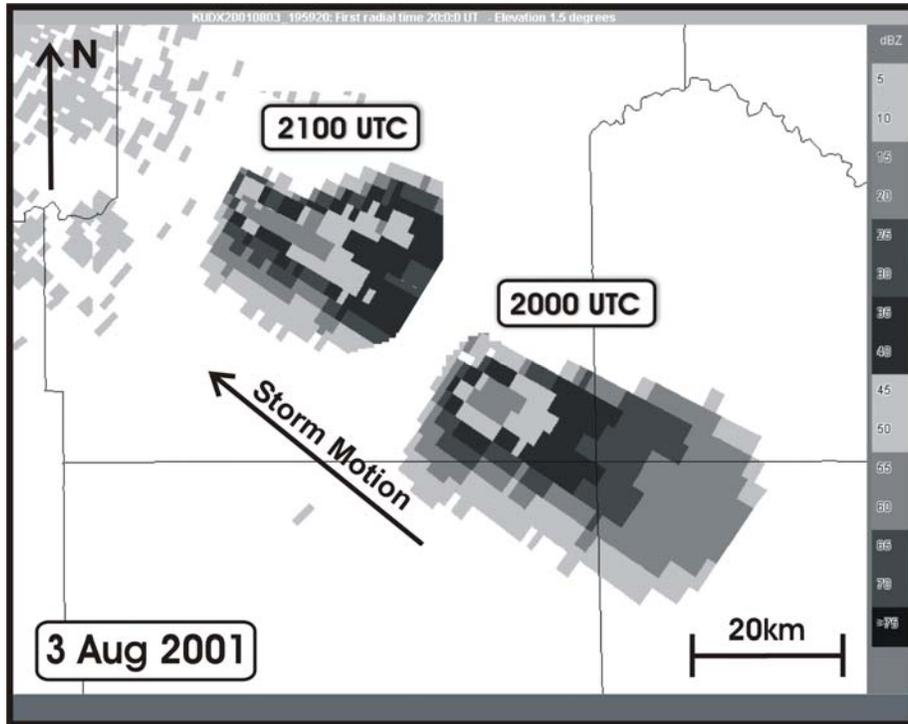
**Fig. 12.** Same as Fig. 9, except for the 0014 UTC 31 May 1999 KHGX VWP hodograph.



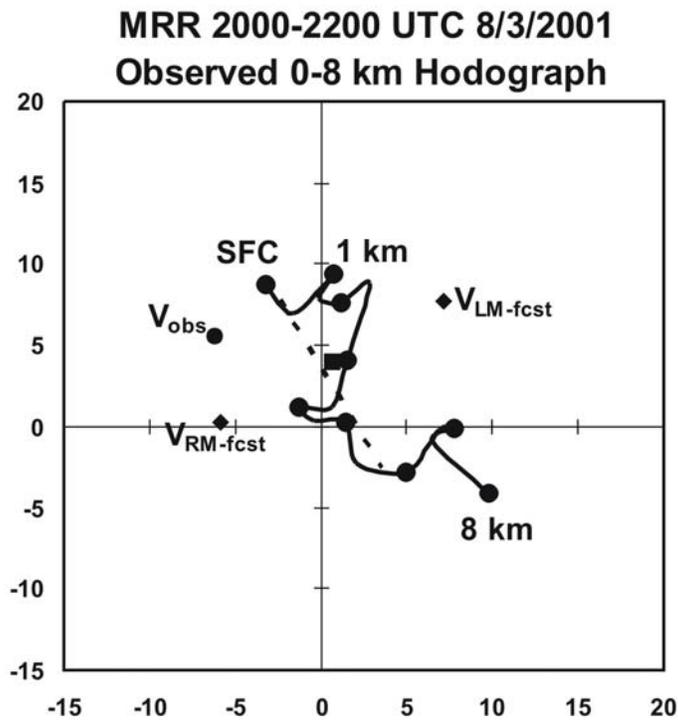
**Fig. 13.** Map of left-moving (LM) and right-moving (RM) storm locations and paths from 2303 UTC 26 March 2000 to 0110 UTC 27 March 2000. KAUS is Austin-Bergstrom International Airport; KCLL is Easterwood Field-College Station; KHGX is the Houston/Galveston WSR-88D radar; KSAT is San Antonio International Airport. The Java-IRAS software was used to display the radar data (Preignitz 1995).



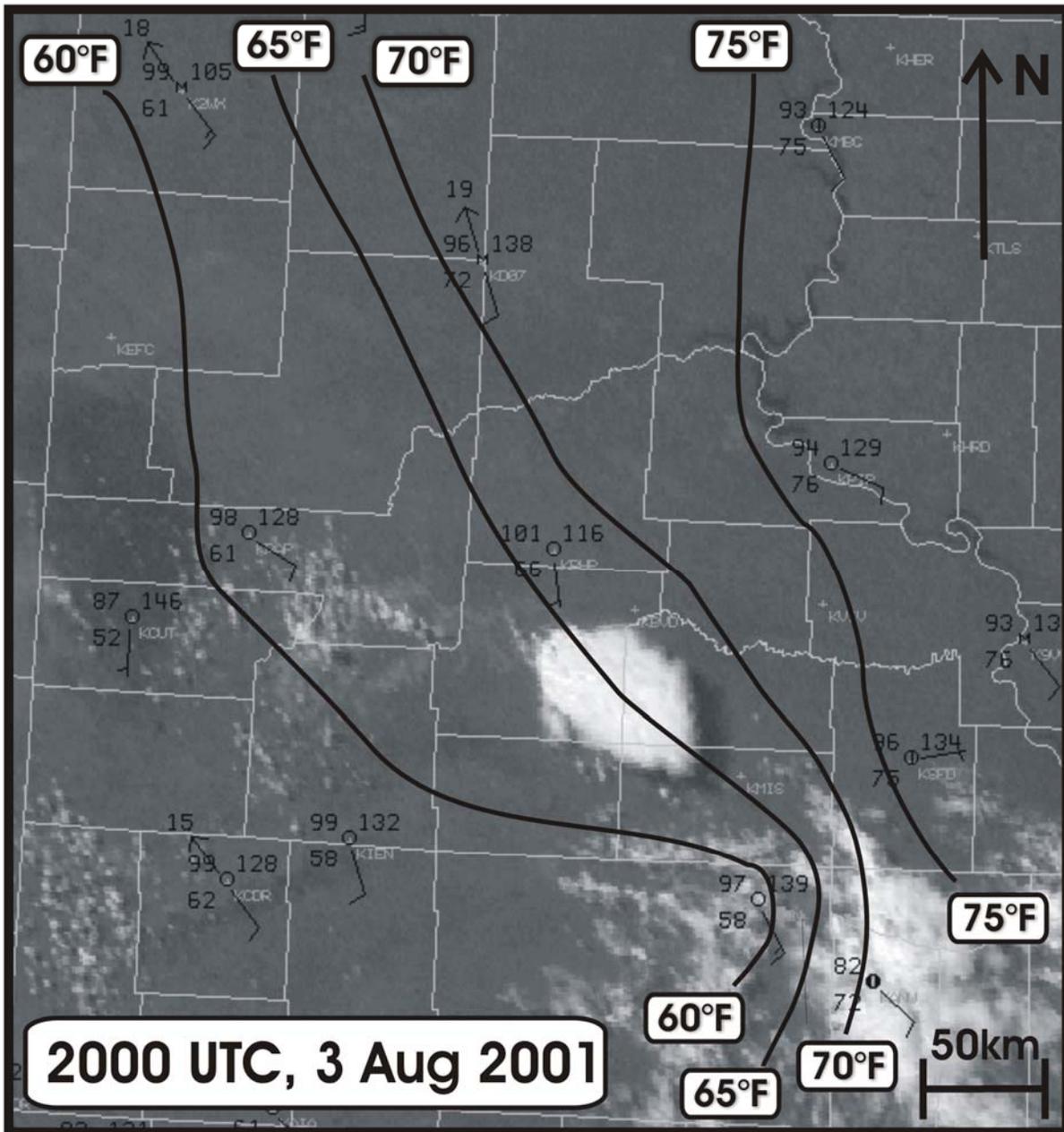
**Fig. 14.** Same as Fig. 9, except for the 0100 UTC 3/27/2000 MAPS (KGRK) hodograph, and  $V_{obs}$  for the right-moving and left-moving supercells. Scale is also different from Figs. 9-12.



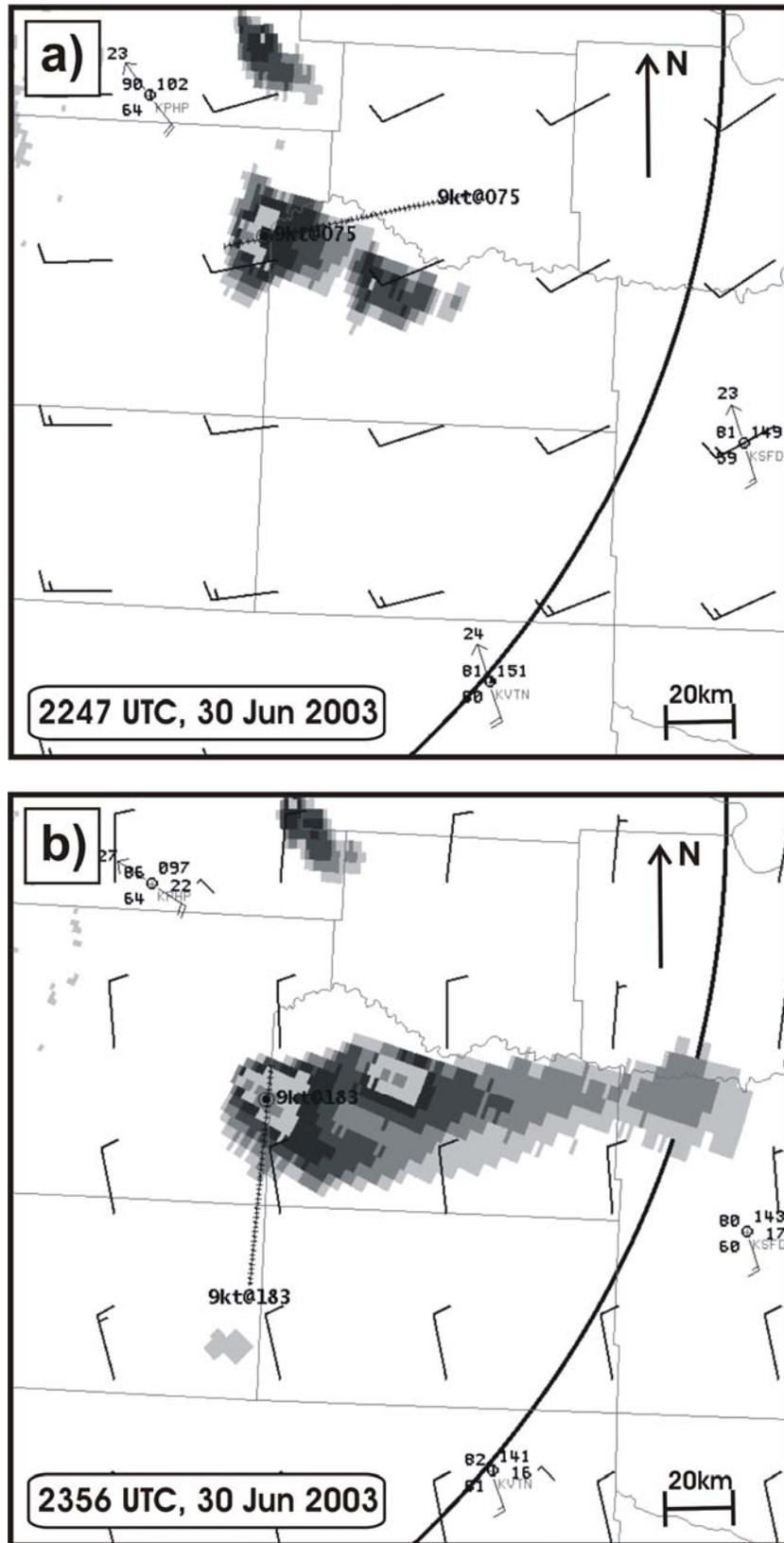
**Fig. 15.** KUDX WSR-88D 1.5° reflectivity composite for 2000 UTC and 2100 UTC 3 August 2001. The storm location at 2000 UTC corresponds to the satellite image in Fig. 17. The Java-IRAS software was used to display the radar data (Preignitz 1995).



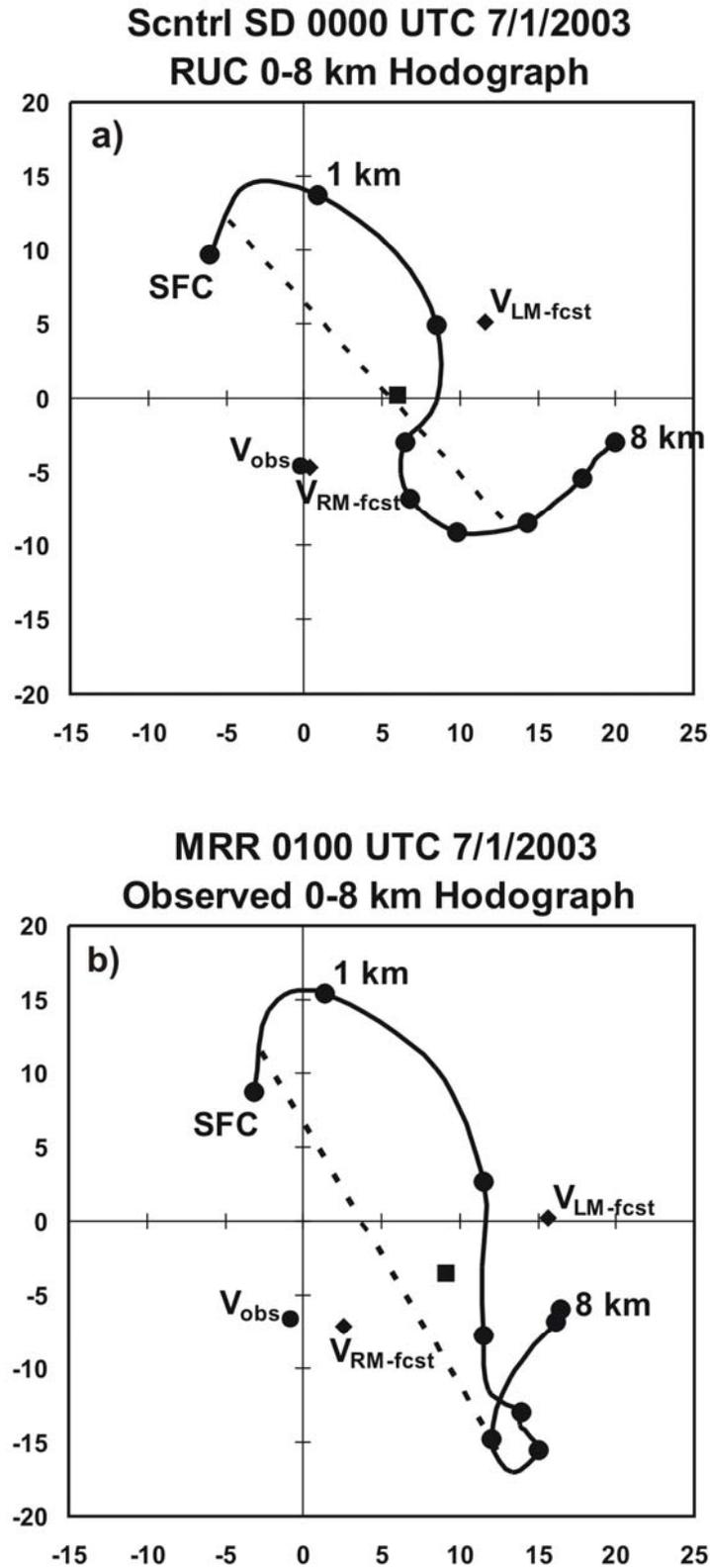
**Fig. 16.** Same as Fig. 9, except for the 2000-2200 UTC 3 August 2001 Merriman, NE, profiler hodograph.



**Fig. 17.** Visible satellite image at 2000 UTC 3 August 2001. Surface observations are overlaid in black, and subjectively drawn isodrosotherms are given for 60, 65, 70, and 75° F. For reference, the Black Hills are located in the center of the left-hand side of the image.



**Fig. 18.** KUDX WSR-88D 0.5° reflectivity image at (a) 2247 UTC 30 June 2003 and (b) 2356 UTC 30 June 2003. Reflectivity shading corresponds to that in Fig. 15. The observed storm motion for the previous 45-min period is displayed with respect to the direction the storm is traveling (knots). Surface observations are overlaid in black. The 0-6 km mean wind vectors (knots) from the RUC model are plotted on (a). The predicted supercell motion vectors from the RUC model using the method from Bunkers et al. (2000) are plotted on (b).



**Fig. 19.** Same as Fig. 9, except for (a) the 0000 UTC 1 July 2003 RUC hodograph and (b) the 0100 UTC 1 July 2003 Merriman, NE, profiler hodograph. Scale is also different from Figs. 9-12.